

# Molecular synthesis in hypervelocity impact plasmas on the primitive Earth and in interstellar clouds

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[1] When impact speeds exceed a critical range of  $\sim 15\text{--}20\text{ km s}^{-1}$ , the resulting plasma plume energizes very rapidly from a few eV, to as high as  $30\text{--}50\text{ eV}$  when speeds reach  $50\text{--}100\text{ km s}^{-1}$ , and can be completely atomized and ionized. During post-impact adiabatic dispersion, the degree of ionization drops. The remaining ion population is mostly atomic, sampling the elemental composition of the colliding bodies, but some molecular ions are also formed. We present here results of experiments in which the physical and chemical conditions in the plume of a hypervelocity micrometeorite impact well above  $20\text{ km s}^{-1}$  were modeled with Q-switched laser ablation. The formed molecular ions were analyzed with a time-of-flight (TOF) mass reflectron capable of distinguishing plasma-synthesized from surface-desorbed species. Singly- and multiply-charged small organic molecules were identified as having formed in laser plasmas induced from inorganic carbonaceous substances. The results show the possibility of abiogenic synthesis of organic molecules in hypervelocity impacts of meteors on the primitive Earth during heavy bombardment and of dust particles in interstellar clouds.

INDEX TERMS:

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## 1. Introduction

[2] Hypervelocity impacts are part of the fundamental dynamics of matter in space, occurring statistically over vast spatial and temporal scales [Spitzer, 1978; Bochkarev, 1992]. At the multiple-kilometer scale, catastrophic impacts between small bodies and planets in the early solar system strongly influenced the geology and geochemistry of those bodies. During heavy bombardment,  $4.5\text{--}3.8\text{ Ga}$ , the primitive Earth's surface endured a veritable rain [Goldsmith and Owen, 1992] of hypervelocity meteor impacts at fluxes of  $\sim 10^3\text{ J m}^{-2}\text{ year}^{-1}$  ( $10^{22}\text{ eV m}^{-2}\text{ year}^{-1}$ ), comparable to the heat from volcanoes [Kobayashi and Saito, 2000]. Velocities ranged from  $\sim 10\text{ km s}^{-1}$  to at least  $70\text{ km s}^{-1}$ . Near the end of this period, life arose.

[3] In interstellar space and protosolar nebulae, various mechanisms can accelerate  $\sim 10\text{--}100\text{ nm}$  dust particles to high velocities. For example, the pressure of light or electromagnetic shocks can accelerate dust to well over  $100\text{ km s}^{-1}$  [Bochkarev, 1992] causing "supercritical" impacts with atomization and ionization fractions approaching unity. Laboratory investigations of impacts are typically carried out in special dust accelerators [Knabe and Krüger, 1982; Ratcliff *et al.*, 1997; Roybal *et al.*, 1995], with particles of  $10^{-11}\text{--}10^{-17}\text{ g}$  mass now achieving speeds up to  $100\text{ km s}^{-1}$ . Studies with such apparatus are challenging in terms of throughput and expense, and are somewhat limited in scale and compatibility with real-time composition measurements. "Full-scale" in-space experiments, such as the dust impact spectrometers on the VEGA and Giotto missions, which analyzed the  $\sim 70\text{--}80\text{ km s}^{-1}$  dust at comet Halley [Dallman *et al.*, 1977], are direct probes of composition [Kissel *et al.*, 1986], but such opportunities are rare.

## 2. Experiment

### 2.1. Laser Simulation of Impact Plasma

[4] Pulsed laser ablation is an accessible and flexible alternative for studying impact plasmas [Pirri, 1977; Roybal *et al.*, 1995]. Irradiances up to  $\varepsilon \sim 10^{13}\text{ W cm}^{-2}$  are readily achieved by focusing a  $10\text{--}100\text{ millijoule-class}$  Q-switched laser (e.g., Nd:YAG) to a diameter  $d \sim 10\text{--}10^3\text{ }\mu\text{m}$ . The few-ns pulse duration is comparable to the impact interaction time of a roughly spherical micrometeorite with a diameter of tens of  $\mu\text{m}$  impacting at tens of  $\text{km s}^{-1}$ . However, one must take care when interpreting the laser irradiance  $\varepsilon = \Delta E/At$ , with pulse energy  $\Delta E$ , spot area  $A$ , and duration  $t$ , in terms of the kinematic impact irradiance  $\varepsilon_{\text{HVI}} = KE/\pi r^2 t \sim \rho v^3/3$  associated with a body of density  $\rho$  and velocity  $v$ . A laser pulse couples quite differently to the target than does a solid impactor [Kostin *et al.*, 1997]: its energy is absorbed much more efficiently and over a shallower region. Furthermore, much of the laser pulse is absorbed above the surface in the expanding plume, rapidly energizing volatilized atoms. As such, crater formation, solid shock, and ejecta patterns are qualitatively different between the laser and impact cases. However, the ionizing impact vapor plume itself, during expansion, cooling, and recombination [Hornung *et al.*, 1996], may be effectively simulated with a pulsed laser at irradiances two to four

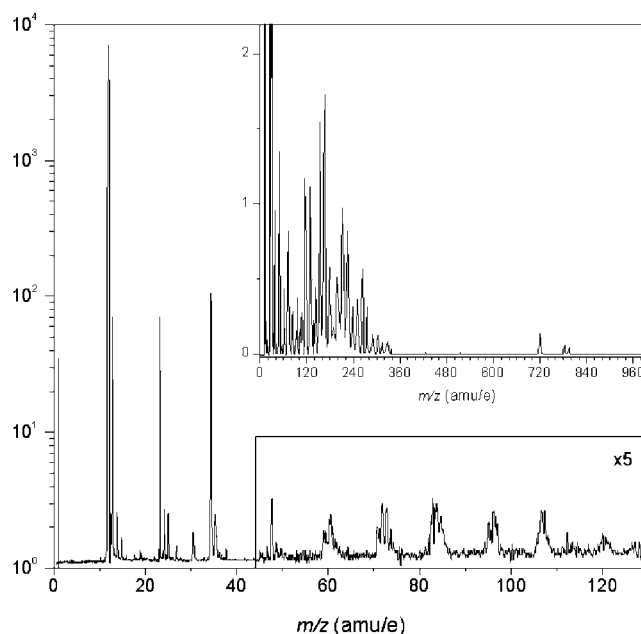
orders of magnitude lower than the “equivalent” micro-meteorite impact irradiance.

[5] The choice of laser parameters for our impact plasma study was based on the measurements by the PUMA instruments on the VEGA spacecraft [Kissel *et al.*, 1986], in which  $\sim 10$  nm to  $1\text{ }\mu\text{m}$  dust particles (mass  $10^{-18}$  to  $10^{-12}$  g) impacted a silver plate at  $\sim 80\text{ km s}^{-1}$ , producing only singly-charged plasma ions with kinetic energies (KEs) of 50–70 eV. The irradiances we used to produce the corresponding laser plasma were in the range of  $\varepsilon \sim 10^9$ – $10^{10}\text{ W cm}^{-2}$ , and estimates have shown that a  $50\text{ }\mu\text{m}$  laser spot size generates a plasma volume comparable to that of a  $10\text{ }\mu\text{m}$  dust particle impact.

[6] The reflectron TOF analyzer used was based upon a technique developed for in situ chemical analysis on planetary probes [Brinckerhoff *et al.*, 2000; Managadze and Shutyayev, 1993]. A Q-switched Nd:YAG laser, with infrared wavelength 1064 nm, maximum pulse energy 25 mJ, and pulse duration 6 ns, was focused through a variable attenuator to a 30–100  $\mu\text{m}$  diameter spot on the target surface *in vacuo* ( $0.1\text{ }\mu\text{Torr}$ ). In contrast to other laser mass spectrometers, ions were not extracted from the source through a large electric potential. The plume formed and expanded in zero external field, as is the case with impacts. A two-stage band-pass reflectron focused a selected KE range onto the detector. As ions formed in different ways attain distinct KE distributions, we were able to distinguish molecules that recombined from atomic ions in the plume (at higher KEs  $\sim 30$  eV) from those desorbed intact from the target (at lower KEs  $\sim 1$  eV). The use of real-time (shot-by-shot) TOF-MS also permitted a “direct” analysis of plasma synthesis [Managadze, 2001a] as opposed to an “indirect” process wherein ion signals from laser pulses are integrated over time. The latter method concentrates trace constituents but does not isolate plume-formed from surface-desorbed species.

## 2.2. Vapor Plume Analyses

[7] Carbon-based target materials were selected for that element’s central role in complex molecular synthesis [Kobayashi and Saito, 2000], and given the small volume of the simulated plume. High purity carbon, with a volumetric component of hydrogen verified by a surface stripping method, and C-rich substances were used. In the mass spectra from pure carbon (Figure 1), singly-ionized, alkanes, alkenes, and other  $\text{C}_N\text{H}_M$  oligomers were identified with  $N = 1$  to 40 (up to  $N = 10$  shown in Figure 1) and  $M = 1$  to 4. Spectra were highly reproducible, although the mass resolution was degraded at the higher masses due to reflectron defocusing in the lower kinetic energy range.  $\text{C}_{60}$  was also detected ( $m/z$  720, Figure 1 inset, averaged signal) in some spectra at slightly higher irradiance. The  $\text{C}_{60}$  signal level was quite low, relative to the smaller masses, when compared to that seen in other studies of the *in vacuo* laser vaporization of graphite at similar irradiances [e.g., Zhang *et al.*, 1999]. We suggest that the limited cluster sizes and the low  $\text{C}_{60}$  yield observed reflect our experimental isolation of the prompt, plasma-synthesized molecular ions from the total population of evaporated species (mainly neutral molecular fragments). Vaporized neutral fragments may serve as precursors to larger clusters which are not detected here. The peaks between  $m/z$  780 and 800 in Figure 1 (inset) were not identified. Spectra similar to Figure 1

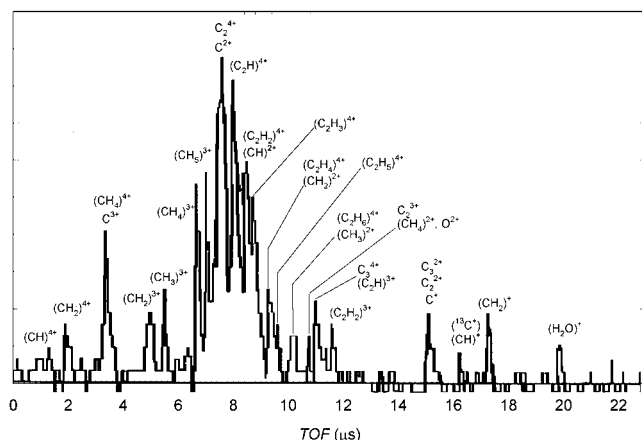


**Figure 1.** Laser mass spectral data for a carbon target, within the kinetic energy range 30–40 eV. The 1064 nm Nd:YAG laser irradiance was about  $2\text{ GW cm}^{-2}$  (pulse duration 6 ns, focus diameter  $50\text{ }\mu\text{m}$ ). Vertical axis is arbitrary logarithmic scale. Inset: Average of mass spectra for irradiances  $\geq 2\text{ GW cm}^{-2}$ .

were also obtained with polycrystalline graphite, paraffin, and bitumen targets [Managadze, 2001a; 2001b].

[8] As the irradiance was increased from  $\varepsilon \sim 1\text{ GW cm}^{-2}$  toward  $10\text{ GW cm}^{-2}$ , multiply-charged  $\text{C}^{2+}$ ,  $\text{C}^{3+}$ , and  $\text{C}^{4+}$  atomic peaks emerged uniformly as expected. However, ablation of the paraffin target at  $\varepsilon \sim 5$ – $10\text{ GW cm}^{-2}$  additionally produced peaks corresponding to multiply-charged hydrocarbon ions (Figure 2). Identifiable peaks were mostly limited to those with  $N = 1$  or 2, but included several where the charge state exceeded  $M$  (e.g.,  $\text{CH}_2^{3+}$  at  $m/z$  4.67), indicating that such ions are not formed only by protonation of neutrals. The formation of such ions in laser induced plasmas does not appear to have been reported previously, but their appearance may be related to multiple charge exchange or stripping processes [Andrews *et al.*, 1992].

[9] A number of factors support the assertion that the observed spectra represent chemical synthesis in the plume, rather than survival (cluster desorption) of molecular fragments [cf. Ramendik *et al.*, 1979; Bulgakov *et al.*, 2000]. First, the high-density plasma produced from carbonaceous targets at these irradiances atomizes its constituents with very high efficiency [Avrorin *et al.*, 1996]. Second, the experimental configuration essentially excludes the detection of ions produced outside the plasma. Without an extraction field, ions formed from desorbed species (such as those from the lower-irradiance region around the central laser spot or from slower thermal processes) are not imparted with sufficiently high kinetic energies to be seen in the analytical window. Third, the multiply-charged hydrocarbons (e.g.,  $\text{CH}_4^{4+}$ ) observed above  $5\text{ GW cm}^{-2}$  occurred simultaneously with singly-charged ions and





energetic impacts on the primitive Earth, potentially “seed-ing” the surface and atmosphere with precursors for processing by longer-duration mechanisms.

[15] Impacts among interstellar dust particles would not generate the large plasma volumes of meteor or small body impacts, but their “universal” prevalence may embody an altogether more substantial source of plasma-synthesized molecular precursors for chemical evolution. Over 100 interstellar molecular species have been detected in astronomical observations to date. These are synthesized from the major biogenic elements (H, C, O, and N) in cold gas phase and shock chemistry, as well as during adsorption of simple precursors on sub-micron ice-mantled dust grains. It is also possible that hypervelocity collisions among dust particles contribute significantly. Assuming a 10 nm interstellar dust particle impact velocity  $v_d \sim 10^2 \text{ km s}^{-1}$ , our data would suggest an atomic plasma mass of  $\sim 3 \cdot 10^{-17} \text{ g}$  per collision, yielding  $N_d \sim 10^3$  molecules of average mass 30 amu, at the observed 0.1% relative yield. Assuming the  $100 \text{ km s}^{-1}$  particles comprise only 10% of the dust population [Spitzer, 1978], the synthetic rate would be modulated by a collision time  $\tau_d \sim [v_d(0.1 n_d)(2\sigma_d)]^{-1} = 5 \cdot 10^{13} \text{ s}$ , where  $n_d \approx 10^{-8} \text{ cm}^{-3}$  and  $\sigma_d \approx 10^{-12} \text{ cm}^2$  are the particle concentration and cross section, respectively [Bochkarev, 1992]. By comparison, the adsorption of hydrogen atoms striking dust particles at  $v_H \sim 1 \text{ km s}^{-1}$  may produce  $\text{H}_2$  at a rate limited by  $\tau'_H \sim (v_H n_d \sigma_d)^{-1} = 10^{15} \text{ s}$ . In a given volume of space, the much higher hydrogen density  $n_H \sim 1 \text{ cm}^{-3} = 10^9 (0.1 n_d)$  leads to a relative collision rate for H of  $\sim 10^9 \tau_d / \tau'_H = 5 \cdot 10^7$  for each such hypervelocity dust-dust impact that produces  $N_d$  molecules.  $\text{H}_2$  formation, even if not highly efficient, must be several orders of magnitude more effective than impact synthesis of molecules. However, as the abundances of C, N, O, and S in the interstellar medium are 3–4 orders of magnitude lower than that of hydrogen, it is possible that the quantity of heavy molecules formed by plasma synthesis is not much less than that formed via surface reactions.

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